# Complete equational theories for the sum-over-paths with unbalanced amplitudes

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Quantum Physics and Logic Paris, July 19, 2023 In the path integral formulation of QM, state evolution is viewed as a sum over the paths taken in superposition

$$|\Psi|i
angle = \sum_{\pi\in\Pi}\psi(\pi)|f(\pi)
angle$$

With certain representations of the sum-over-paths,

- supports efficient composition
- can be re-written symbolically to reduce interfering paths
- highly amenable to automated simplification, verification & simulation of some circuits

## Symbolic representation of balanced sums

Recall: A **balanced** path sum is a symbolic expression of a linear operator  $\Psi : \mathbb{C}^{2^m} \to \mathbb{C}^{2^n}$  as a sum of the form

$$|\Psi|ec{x}
angle = |\Psi(ec{x})
angle = \mathcal{N}\sum_{ec{y}\in\mathbb{Z}_2^k}e^{2\pi i P(ec{x},ec{y})}|f(ec{x},ec{y})
angle,$$

N ∈ C \ {0} is a normalization factor,
 P : Z<sub>2</sub><sup>m</sup> × Z<sub>2</sub><sup>k</sup> → R is a real-valued multilinear polynomial, and
 f : Z<sub>2</sub><sup>m</sup> × Z<sub>2</sub><sup>k</sup> → Z<sub>2</sub><sup>n</sup> is system of n Boolean polynomials

#### Notation

 $|\Psi(x, y, ...)\rangle$  refers to a path sum expression with distinguished free variables x, y, ...

## Examples

Phase & reversible gates:

$$S|x\rangle = i^{x}|x\rangle, \qquad T|x\rangle = e^{rac{i\pi}{4}x}|x
angle$$
  
 $X|x\rangle = |1 \oplus x\rangle \qquad CNOT|x
angle|y
angle = |x
angle|x \oplus y
angle$ 

Branching gates:

$$|H|x
angle = rac{1}{\sqrt{2}}\sum_{y}(-1)^{xy}|y
angle$$

Cups and caps:

$$C = \sum_{y} |y\rangle |y
angle, \qquad \supset |x_1
angle |x_2
angle = rac{1}{2} \sum_{y} (-1)^{y(x_1+x_2)}$$

If  $|\Psi\rangle$  is a path sum where  $y \notin FV(\Psi)$  and f is a Boolean expression such that  $x, y \notin FV(f)$ , then the following hold:

$$\sum_{\substack{y \in \mathbb{Z}_2 \\ x, y \in \mathbb{Z}_2}} |\Psi\rangle \equiv_{\text{Cliff}} 2|\Psi\rangle \tag{E}$$
$$\sum_{\substack{x, y \in \mathbb{Z}_2 \\ y \in \mathbb{Z}_2}} (-1)^{y(x+f)} |\Psi(x)\rangle \equiv_{\text{Cliff}} 2|\Psi(f)\rangle \tag{H}$$
$$\sum_{\substack{y \in \mathbb{Z}_2 \\ y \in \mathbb{Z}_2}} i^y (-1)^{yf} |\Psi\rangle \equiv_{\text{Cliff}} \omega \sqrt{2} (-i)^{\overline{f}} |\Psi\rangle \tag{\omega}$$

Theorem

 $\equiv_{Cliff}$  is complete for Stabilizer operations

## Completeness for Toffoli+Hadamard

$$\begin{split} \sum_{y \in \mathbb{Z}_2} \sum_{x \in \mathbb{Z}_2} (-1)^{y(x \cdot g + g \cdot f + 1)} |\Psi(x)\rangle &\equiv_{\mathrm{TH}} \sum_{y \in \mathbb{Z}_2} (-1)^{y(g+1)} |\Psi(1+f)\rangle \quad (\mathsf{Hgen}) \\ &\sum_{y \in \mathbb{Z}_2} \sum_{x \in \mathbb{Z}_2} (-1)^{y \cdot f + x \cdot g} |\Psi\rangle \equiv_{\mathrm{TH}} 2 \sum_{y \in \mathbb{Z}_2} (-1)^{y(f+g+f \cdot g)} |\Psi\rangle \quad (\mathsf{Hrel}) \\ &\sum_{y \in \mathbb{Z}_2} \sum_{x \in \mathbb{Z}_2} (-1)^y |\Psi\rangle \equiv_{\mathrm{TH}} 0 |\Psi\rangle \quad (\mathsf{Z}) \end{split}$$

#### Theorem (Vilmart 2022)

 $\equiv_{TH+Cliff}$  is complete for Toffoli-Hadamard and by extension Clifford+ $R_Z(2\pi/2^k)$ 

Renaud Vilmart, Completeness of Sum-Over-Paths for Toffoli-Hadamard and the Dyadic Fragments of Quantum Computation. CSL 2023.

Goal: can we translate the full ZH-calculus to a complete equational theory for arbitrary path sums?

# From ZH to SOP

#### Recall: The ZH-calculus<sup>1</sup> is generated by

$$\llbracket n ] = \sum_{x} |xx \cdots x\rangle \langle xx \cdots x|$$
$$\llbracket n ] = \sum_{x, \vec{y}} \alpha^{x_1 \cdots x_n y_1 \cdots y_m} |\vec{y}\rangle \langle \vec{x}|$$

<sup>&</sup>lt;sup>1</sup>Miriam Backens and Aleks Kissinger, ZH: A Complete Graphical Calculus for Quantum Computations Involving Classical Non-linearity. QPL 2018.

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With a slight tweak we can write them in the style of path sums:

$$\begin{bmatrix} n \\ \vdots \\ \infty \\ \vdots \\ m \end{bmatrix} |\vec{x}\rangle = \sum_{y} \sum_{\vec{z}} 2^{-n} (-1)^{\sum_{i=1}^{n} z_i(x_i+y)} |yy \cdots y\rangle$$
$$\begin{bmatrix} n \\ \vdots \\ \infty \\ \cdots \\ m \end{bmatrix} |\vec{x}\rangle = \sum_{\vec{y}} \alpha^{x_1 \cdots x_n y_1 \cdots y_m} |\vec{y}\rangle$$

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#### Problem: need a notion of unbalanced sums!

<sup>1</sup>Miriam Backens and Aleks Kissinger, ZH: A Complete Graphical Calculus for Quantum Computations Involving Classical Non-linearity. QPL 2018.

#### Balanced vs Unbalanced

The controlled-H gate

$${\cal CH} = rac{1}{\sqrt{2}} egin{bmatrix} \sqrt{2} & 0 & 0 & 0 \ 0 & \sqrt{2} & 0 & 0 \ 0 & 0 & 1 & 1 \ 0 & 0 & 1 & -1 \end{bmatrix}$$

is unbalanced since non-zero entries differ in magnitude

Balanced representations of unbalanced operators like the controlled -H gate are tricky to design and simplify!

$$CH|x_1x_2\rangle = rac{1}{\sqrt{2}}\sum_{y\in\mathbb{Z}_2}\omega^{(1-x_1)(2y-1)}(-1)^{x_1x_2y}|x_1
angle|(1\oplus x_1)x_2\oplus x_1y
angle$$

An unbalanced sum-over-paths is an expression  $|\Psi\rangle$  of the following language

$$f ::= 0 \mid 1 \mid x, y, z, \dots \mid f_1 \cdot f_2 \mid f_1 \oplus f_2 \mid \neg f := 1 \oplus f$$
$$a ::= \alpha, \beta \in \mathcal{F} \mid a^f \mid a_1 a_2$$
$$|\Psi\rangle ::= \sum_{\vec{y}} a \mid f_1 \cdots f_n \rangle.$$

where variables  $x, y, z, \ldots$  range over Booleans and  $\mathcal{F}$  is a field.

## Universality

Unbalanced sums subsume balanced sums and are trivially universal as any linear operator  $\mathcal{F}^{2^n} \to \mathcal{F}^{2^m}$  may be uniquely written as

$$\Psi |\vec{x}\rangle = \sum_{\vec{y}} \alpha_{00\cdots0}^{00\cdots0=\vec{x}\vec{y}} \alpha_{00\cdots1}^{00\cdots1=\vec{x}\vec{y}} \cdots \alpha_{11\cdots1}^{11\cdots1=\vec{x}\vec{y}} |\vec{y}\rangle$$

where  $\vec{x} = \vec{y} := \prod_i x_i \oplus y_i \oplus 1$ .

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where  $\vec{x} = \vec{y} := \prod_i x_i \oplus y_i \oplus 1$ .

#### Definition (Normal form)

A normal form is a closed, unbalanced sum of the following form:

$$\sum_{\vec{x}} \alpha_{00\cdots0}^{00\cdots0=\vec{x}} \alpha_{00\cdots1}^{00\cdots1=\vec{x}} \cdots \alpha_{11\cdots1}^{11\cdots1=\vec{x}} |\vec{x}\rangle$$

#### Normal forms Unbalanced sums vs. ZH calculus

#### Unbalanced sums:

$$\sum_{\vec{x}} \alpha_{00\cdots0}^{00\cdots0=\vec{x}} \alpha_{00\cdots1}^{00\cdots1=\vec{x}} \cdots \alpha_{11\cdots1}^{11\cdots1=\vec{x}} |\vec{x}\rangle$$





#### Equational theory of the ZH-calculus

(ZS1)



 $\left\{ \begin{array}{c} \\ \\ \\ \end{array} \right\} = \left\{ \begin{array}{c} \\ \\ \end{array} \right\} = \left\{ \begin{array}{c} \\ \\ \end{array} \right\}$ 









(M)





= 6 (U) Н

(A)







## Equational theory of the ZH-calculus

 $\begin{array}{c} & & \\$ (ZS1) (A) 2  $\frac{a+b}{2}$ (I)  $\left| \begin{array}{c} \\ \\ \\ \\ \\ \end{array} \right| = \left| \begin{array}{c} \\ \\ \\ \\ \end{array} \right| = \left| \begin{array}{c} \\ \\ \\ \\ \end{array} \right|$ a o  $(HS2) \qquad = 2$ a à (ZS2) (BA2) (O) 2 = (...) = (BA1) = m  $\downarrow$  =  $\downarrow$ (M) (U) ab

Only need analogues of (A) (I) and (O)

## Formulating the average rule

ZH-calculus:



# Formulating the average rule

ZH-calculus:



Unbalanced sums:

$$\sum_{y} (\alpha^{y} \beta^{\neg y})^{x} |\Psi\rangle = 2 \left(\frac{\alpha + \beta}{2}\right)^{x} |\Psi\rangle$$

#### Formulating the average rule

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Unbalanced sums:

$$\sum_{y} (\alpha^{y} \beta^{\neg y})^{x} |\Psi\rangle = 2 \left(\frac{\alpha + \beta}{2}\right)^{x} |\Psi\rangle$$

In SOP-land, we can generalize x to arbitrary expressions:

$$\sum_{y} \left( \alpha^{y} \beta^{\neg y} \right)^{f} |\Psi\rangle \equiv 2 \left( \frac{\alpha + \beta}{2} \right)^{f} |\Psi\rangle$$

# Formulating the intro rule

ZH-calculus:



# Formulating the intro rule

ZH-calculus:



Unbalanced sums:

$$\sum_{y} a^{x} |y\rangle = \sum_{y} a^{xy} a^{x(\neg y)} |y\rangle$$

#### Formulating the intro rule

ZH-calculus:



Unbalanced sums:

$$\sum_{y} a^{x} |y
angle = \sum_{y} a^{xy} a^{x(\neg y)} |y
angle$$

Generalized rule:

$$a \equiv a^{y} a^{\neg y}$$

# Formulating the ortho rule

ZH-calculus:



# Formulating the ortho rule

ZH-calculus:



Unbalanced sums:

$$2\sum_{w,y,z}(-1)^{(\neg x)wy}(-1)^{xwz}|yz\rangle = \sum_{w_1,w_2,y,z}(-1)^{(\neg x)w_1y}(-1)^{xw_2z}|yz\rangle$$

#### Formulating the ortho rule

ZH-calculus:



Unbalanced sums:

$$2\sum_{w,y,z}(-1)^{(\neg x)wy}(-1)^{xwz}|yz\rangle = \sum_{w_1,w_2,y,z}(-1)^{(\neg x)w_1y}(-1)^{xw_2z}|yz\rangle$$

Thus a rule is born:

$$\sum_{y,z} a_1(y)^x a_2(z)^{\neg x} |\Psi(x)\rangle \equiv 2 \sum_y a_1(y)^x a_2(y)^{\neg x} |\Psi(x)\rangle$$

## ZH vs. unbalanced sums



The rest

unbalanced sums  $\sum_{\vec{x}} \alpha_{00...0}^{00...0=\vec{x}} \alpha_{00...1}^{00...1=\vec{x}} \cdots \alpha_{11...1}^{11...1=\vec{x}} |\vec{x}\rangle$  $\sum_{y} (\alpha^{y} \beta^{\neg y})^{f} |\Psi\rangle \equiv 2 \left(\frac{\alpha + \beta}{2}\right)^{f} |\Psi\rangle$  $a \equiv a^f a^{\neg f}$  $\sum_{y,z} a_1(y)^x a_2(z)^{\neg x} |\Psi(x)\rangle \equiv 2\sum_y a_1(y)^x a_2(y)^{\neg x} |\Psi(x)\rangle$ 

 $\sum_{x,y}(-1)^{y(x\oplus f)}|\Psi(x)
angle\equiv 2|\Psi(f)
angle$ 

#### Rules of the unbalanced sum-over-paths

$$f \oplus 0 \equiv_{\mathcal{F}} f \qquad (f_1 \oplus f_2) \oplus f_3 \equiv_{\mathcal{F}} f_1 \oplus (f_2 \oplus f_3)$$

$$f \oplus f \equiv_{\mathcal{F}} 0 \qquad f_1 \oplus f_2 \equiv_{\mathcal{F}} f_2 \oplus f_1$$

$$f \cdot 1 \equiv_{\mathcal{F}} f \qquad (f_1 \cdot f_2) \cdot f_3 \equiv_{\mathcal{F}} f_1 \cdot (f_2 \cdot f_3)$$

$$f \cdot f \equiv_{\mathcal{F}} f \qquad f_1 \cdot f_2 \equiv_{\mathcal{F}} f_2 \cdot f_1$$

$$f_1 \cdot (f_2 \oplus f_3) \equiv_{\mathcal{F}} f_1 \cdot f_2 \oplus f_1 \cdot f_3$$

$$\begin{array}{ll} (a_1 \cdot a_2) \cdot a_3 \equiv_{\mathcal{F}} a_1 \cdot (a_2 \cdot a_3) \\ a_1 \cdot a_2 \equiv_{\mathcal{F}} a_2 \cdot a_1 \\ a \cdot 1 \equiv_{\mathcal{F}} a \end{array} \qquad \begin{array}{ll} a^0 \equiv_{\mathcal{F}} 1 \equiv_{\mathcal{F}} 1^f \\ a^1 \equiv_{\mathcal{F}} a \equiv_{\mathcal{F}} a^{f_1} a^{-f_1} \end{array} \qquad \begin{array}{ll} a^{f_1 \oplus f_2} \equiv_{\mathcal{F}} a^{f_1} a^{f_2} (a^{-2})^{f_1 \cdot f_2} \\ a^{f_1 \cdot f_2} \equiv_{\mathcal{F}} (a^{f_1})^{f_2} \\ a^{f_1 a_2} \equiv_{\mathcal{F}} (a^{f_1})^{f_2} \end{array} \\ \begin{array}{l} a^{f_1 \oplus f_2} \equiv_{\mathcal{F}} (a^{f_1})^{f_2} \\ a^{f_1 a_2} \equiv_{\mathcal{F}} (a_1 a_2)^f \end{array} \end{array}$$

$$\sum_{x,y} (-1)^{y(x\oplus f)} |\Psi(x)\rangle \equiv_{\mathcal{F}} 2|\Psi(f)\rangle \tag{H}$$

$$\sum_{y,z} a_1(y)^{x} a_2(z)^{\neg x} |\Psi(x)\rangle \equiv_{\mathcal{F}} 2 \sum_{y} a_1(y)^{x} a_2(y)^{\neg x} |\Psi(x)\rangle$$
(0)

$$\sum_{y} (\alpha^{y} \beta^{\neg y})^{f} |\Psi\rangle \equiv_{\mathcal{F}} 2\left(\frac{\alpha + \beta}{2}\right)^{t} |\Psi\rangle \tag{A}$$

Part 1: Normalization of amplitudes

#### Proposition ( $\mathcal{F}$ -expression normalization)

An F-expression a can be brought into the form

$$\alpha_{00\cdots0}^{00\cdots0=\vec{x}}\alpha_{00\cdots1}^{00\cdots1=\vec{x}}\cdots\alpha_{11\cdots1}^{11\cdots1=\vec{x}}.$$

over the variables  $\{x_i\} \supseteq FV(a)$  using the equations of  $\equiv_{\mathcal{F}}$ .

#### Completeness

Part 2: Normalization of sums

#### Theorem

 $\equiv_{\mathcal{F}}$  is complete for unbalanced sums over any field  $\mathcal{F}$ .

Proof sketch:

- 1. Apply (H) to the sum  $\sum_{\vec{y}} a |f_1 f_2 \cdots f_n\rangle$  to get  $\sum_{\vec{x}} \sum_{\vec{y}} a' |\vec{x}\rangle$
- 2. Normalize a' to get

$$\sum_{\vec{x}} \sum_{\vec{y}} \alpha_{00\cdots00}^{00\cdots00=\vec{x}\vec{y}} \alpha_{00\cdots01}^{00\cdots01=\vec{x}\vec{y}} \cdots \alpha_{11\cdots11}^{11\cdots11=\vec{x}\vec{y}} |\vec{x}\rangle$$

3. Sum amplitudes in pairs with (O) and (A) to remove each  $y_i$ :

$$\begin{split} &\sum_{\vec{x}} \sum_{y} (\alpha_{00\cdots00}^{\gamma y} \alpha_{00\cdots01}^{y})^{00\cdots0=\vec{x}} \cdots (\alpha_{11\cdots10}^{\gamma y} \alpha_{11\cdots11}^{y})^{11\cdots1=\vec{x}} |\vec{x}\rangle \\ &\equiv_{\mathcal{F}} \sum_{\vec{y}} \sum_{\vec{y}} \frac{1}{2^{2^{n}-1}} (\alpha_{00\cdots00}^{\gamma y_{1}} \alpha_{00\cdots01}^{y_{1}})^{00\cdots0=\vec{x}} \cdots (\alpha_{11\cdots10}^{\gamma y_{n}} \alpha_{11\cdots11}^{y_{n}})^{11\cdots1=\vec{x}} |\vec{x}\rangle \\ &\equiv_{\mathcal{F}} \sum_{\vec{x}} \frac{2^{2^{n}}}{2^{2^{n}-1}} \left( \frac{\alpha_{00\cdots00} + \alpha_{00\cdots01}}{2} \right)^{00\cdots0=\vec{x}} \cdots \left( \frac{\alpha_{11\cdots10} + \alpha_{11\cdots11}}{2} \right)^{11\cdots1=\vec{x}} |\vec{x}\rangle \end{split}$$

The  $\oplus$ -exponent equation  $a^{f_1 \oplus f_2} \equiv a^{f_1} a^{f_2} (a^{-2})^{f_1 \cdot f_2}$  is not generally admissible over rings, e.g.  $\mathbb{Z}[\frac{1}{\sqrt{2}}, i]$ .

We can instead use  $a^{f_1 \oplus f_2} \equiv a^{f_1} + a^{f_2} - (2a)^{f_1 \cdot f_2}$  if we allow amplitude expressions to be summed...

$$f ::= 0 \mid 1 \mid x, y, z, \dots \mid f_1 \cdot f_2 \mid f_1 \oplus f_2$$
$$r ::= \alpha, \beta \in \mathcal{R} \mid r^f \mid r_1 r_2 \mid r_1 + r_2$$
$$|\Psi\rangle ::= \sum_{\vec{y}} r \mid f_1 \cdots f_n \rangle.$$

where  ${\mathcal R}$  is a commutative ring

#### Normalization of amplitude expressions

$$(r_{1} + r_{2}) + r_{3} \equiv r_{1} + (r_{2} + r_{3})$$

$$r_{1} + r_{2} \equiv r_{2} + r_{1}$$

$$r_{1} + 0 \equiv r_{1}$$

$$r - r \equiv 0$$

$$(r_{1} \cdot r_{2}) \cdot r_{3} \equiv r_{1} \cdot (r_{2} \cdot r_{3})$$

$$r_{1} \cdot r_{2} \equiv r_{2} \cdot r_{1}$$

$$r \cdot 1 \equiv r$$

$$r_{1} \cdot (r_{2} + r_{3}) \equiv r_{1} \cdot r_{2} + r_{1} \cdot r_{3}$$

$$r^{0} \equiv 1 \equiv 1^{f}$$

$$r^{1} \equiv r \equiv r^{f} r^{\neg f}$$

$$r^{f_{1} \oplus f_{2}} \equiv r^{f_{1}} + r^{f_{2}} - (2r)^{f_{1} \cdot f_{2}}$$

$$r^{f_{1} \cdot f_{2}} \equiv (r^{f_{1}})^{f_{2}}$$

$$r^{f}_{1} r^{f}_{2} \equiv (r_{1} r_{2})^{f}$$

$$r^{f}_{1} r^{\neg f}_{2} \equiv r^{f}_{1} + r^{\neg f}_{2} - 1$$

$$r^{f}_{1} + r^{f}_{2} \equiv (r_{1} + r_{2})^{f} + 0^{\neg f}$$

#### Proposition

An amplitude expression r can be brought into the form

$$\alpha_{00\cdots0}^{00\cdots0=\vec{x}}\alpha_{00\cdots1}^{00\cdots1=\vec{x}}\cdots\alpha_{11\cdots1}^{11\cdots1=\vec{x}}$$

over the variables  $\{x_i\} \supseteq FV(r)$  using the above equations (+ axioms of Boolean rings).

$$\sum_{x,y} (-1)^{y(x \oplus f)} |\Psi(x)\rangle \equiv_{\mathcal{R}} 2|\Psi(f)\rangle$$
(H)  
$$\sum_{y} r(y) |\Psi\rangle \equiv_{\mathcal{R}} (r(0) + r(1))|\Psi\rangle$$
(S)

#### Theorem

 $\equiv_{\mathcal{R}}$  is complete for unbalanced sums over any field  $\mathcal{R}$ .

#### Proof sketch:

Use (H) to write the sum as  $\sum_{\vec{y}} \sum_{\vec{z}} r(\vec{y}, \vec{z}) |\vec{z}\rangle$ , then explicitly expand the sum over  $\vec{y}$  with (S) and normalize the amplitude expression

#### More restrictive path sums

A more direct ZH analogue removes sums from both the **Amplitude and Boolean sub-languages**:

$$f ::= 0 \mid 1 \mid x, y, z, \dots \mid f_1 \cdot f_2 \mid f_1 \oplus f_2$$
$$r ::= \alpha, \beta \in \mathcal{R} \mid r^f \mid r_1 r_2 \mid r_1 + r_2$$
$$\mid \Psi \rangle ::= \sum_{\vec{y}} r \mid f_1 \cdots f_n \rangle.$$

In this case, we need rules explicitly handling distribution of Boolean sums, e.g. the following which corresponds to the BA2 rule of the ZH-calculus

$$\sum_{x,y} (-1)^{xy} (-1)^{yf_1} \cdots (-1)^{yf_k} (-1)^{xg} |\Psi\rangle \equiv 2 (-1)^{f_1g} \cdots (-1)^{f_kg} |\Psi
angle,$$

So what did we actually do?

- Defined a notion of unbalanced path sums
- ► Gave complete equational theories over Rings & Fields

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Takeaway...

Symbolic algebra gives a lot of flexibility in the design and operation of formal methods for QC

Thank you!