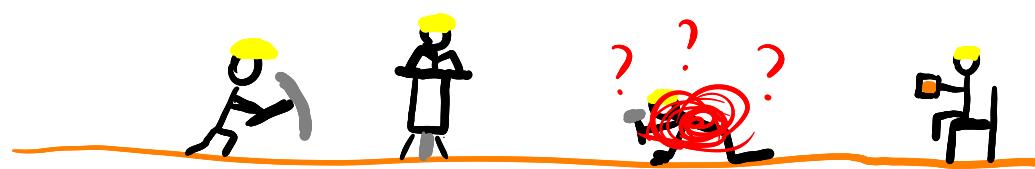


# CMPT 476 Lecture 5

## Working with a qubit



We know now that an *isolated quantum system* corresponds to a  $d$ -dimensional Hilbert Space  $\mathbb{C}^d$  and we can affect its state by applying either:

- Unitary operations  $\xrightarrow{\text{U}}$
- Measurements  $\xrightarrow{\text{M}}$

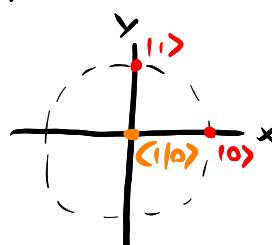
Before we move on to *multiple qubits*, let's see what kind of *quantum effects* we can witness with a *single qubit*.

## (Quantum Zeno effect)

Measurement can effect states in a strange way. For instance, given the state  $|4\rangle = |0\rangle$ , what is the probability of measuring  $|1\rangle$ ?

$$|\langle 1|4\rangle|^2 = 0$$

Geometrically, this is the projection onto the y-axis



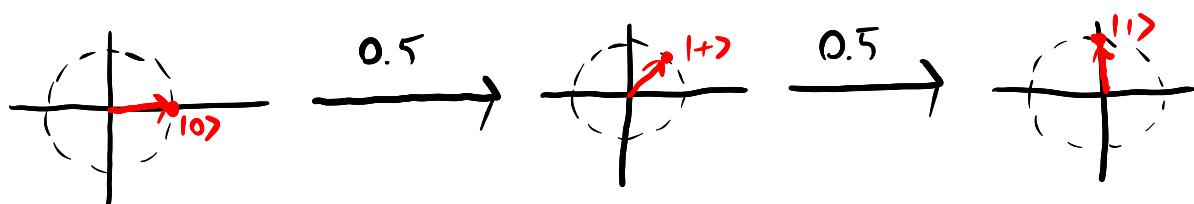
However, suppose we measured first in the  $|+\rangle, |-\rangle$ ? We would get state  $|+\rangle$  with probability

$$\begin{aligned} |\langle +|4\rangle|^2 &= \left| \frac{1}{\sqrt{2}} (\langle 0| + \langle 1|) |0\rangle \right|^2 \\ &= \left| \frac{1}{\sqrt{2}} (\langle 0|0\rangle + \langle 1|0\rangle) \right|^2 \\ &= \frac{1}{2} \end{aligned}$$

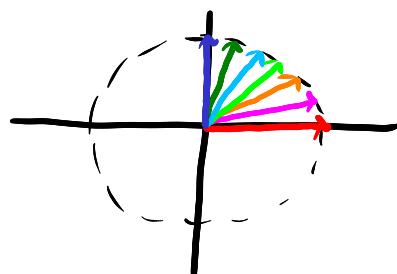
Now when we measure what is the prob. of getting  $|1\rangle$ ?

$$|\langle 1|4\rangle|^2 = \left| \frac{1}{\sqrt{2}} \right|^2 = \frac{1}{2}$$

In effect we can change a state by measuring in different bases



If we make the angle between bases small enough, we can do this with high probability

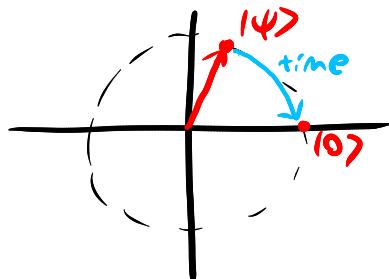


(A watched pot never boils)

A similar effect can be used to control decoherence. Suppose we have a superposition

$$|\Psi\rangle = a|0\rangle + b|1\rangle$$

which decoheres over time to the  $|0\rangle$  state, i.e.



then we can keep it in state  $|\Psi\rangle$  with high probability if we repeatedly measure in the basis

$$\{|\Psi\rangle, |\Psi^\perp\rangle = b^*|0\rangle - a^*|1\rangle\}$$

Of course, this assumes we know  $|\Psi\rangle$  ahead of time, so it's not really an **uncertain state** in a useful sense. We will see that quantum algorithms **require** genuine superpositions where  $a$  and  $b$  are not known *a priori* to achieve speed ups.

## (Elitzur-Vaidman Bomb)

A closely related thought experiment is the Elitzur-Vaidman Bomb. Here are the rules:

A suspicious man hands you and a friend two boxes. With them is a note that reads

"Do you want to play a game? In one box is a bomb triggered by a horizontally polarized photon. If you open the box it will explode, and if you do nothing I will trigger the bomb. Find out which box has the bomb or many will die..."

So, the rules are:

1. We can send a photon in state  $|0\rangle$  into the box
2. If there is no bomb we get state  $|0\rangle$  back
3. If there is a bomb,  $|0\rangle$  is measured
  - 3.1 If the result is  $|0\rangle$  the bomb doesn't go off
  - 3.2 If  $|1\rangle$  the bomb goes off



The thought experiment shows that with high probability a quantum system can win the game.

Suppose our photon **decreases** towards  $|1\rangle$  at an angle of  $\epsilon$  each second.

If we start in state  $|0\rangle$ , and send it into the box **every second**, after  $\frac{\pi}{2\epsilon}$  seconds,

1. If no bomb, state has rotated  $\frac{\pi}{2\epsilon} \cdot \epsilon = 90^\circ$  to  $|1\rangle$
2. If there is a bomb, each time the measurement **sends** us back to the  $|0\rangle$  state if it doesn't trigger, which happens with probability

$$\sin^2 \epsilon \approx \epsilon^2 \text{ for small } \epsilon$$

(remember:  $e^{i\epsilon} = \cos \epsilon + i \sin \epsilon$ )

So we only set off the bomb with probability

$$\frac{\pi}{2\epsilon} \cdot \epsilon^2 = \frac{\pi}{2}\epsilon$$

and if not, we end with state 10 in case of a bomb,  
and 11 if no bomb. Pretty cool!

## (Distinguishing states)

Now that we've had our fun, a more practical question is: given a qubit  $|1\rangle$ , can we determine what  $|1\rangle$  is? It should be fairly obvious that we can't in general if  $|1\rangle$  is unknown, because measuring will collapse the state. In analogy to probabilistic computation, we can't determine the probability distribution of a bit (i.e. its **state**) by observing its value (i.e. **measuring**). However, in some cases we can determine with high probability which of two possible states we have.

### Ex.

Suppose you're handed a qubit  $|1\rangle$  and told that either  $|1\rangle = |0\rangle$  or  $|1\rangle = |+\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$ . Can you determine which case it is?

A simple protocol with a **one-sided error** is to measure in the computational basis and guess

- $|0\rangle$  if the result is 0
- $|+\rangle$  if the result is 1

If  $|1\rangle = |0\rangle$ , then this protocol always guesses correctly. If  $|1\rangle = |+\rangle$ , then we measure 1 with 50% probability and hence guess correctly with 50% probability.

What if the person who gave you  $|1\rangle$  is trying to trick you and intentionally gives  $|+\rangle$  in anticipation of this strategy? In this case it's better to have a **two-sided error**. We can do this by making our guess probabilistically to account for the imbalance.

Here's a two-sided error protocol:

- If result is 0, guess  $|0\rangle$  with probability  $\frac{2}{3}$   
and  $|+\rangle$  with probability  $\frac{1}{3}$
- If result is 1, guess  $|+\rangle$  as before

Now if  $|4\rangle = |0\rangle$ , we guess correctly with prob.  $\frac{2}{3}$  and  
if  $|4\rangle = |+\rangle$ , then we guess correctly with prob.

$$\underbrace{\frac{1}{2} \cdot \frac{1}{3}}_{\text{measure 0}} + \underbrace{\frac{1}{2}}_{\text{measure 1}} = \frac{2}{3}$$

## (Global vs relative phase)

One broad class of states which **cannot be distinguished** are those which differ by a global phase  $e^{i\Theta}$

Ex.

If  $|4\rangle$  is a state, then so is  $|4\rangle - e^{i\Theta}|4\rangle$  for any  $\Theta$ :

$$\langle \psi' | \psi' \rangle = (e^{-i\Theta} \langle 4 |) (e^{i\Theta} | 4 \rangle) = \langle 4 | 4 \rangle$$

$|4\rangle$  and  $|4'\rangle$  are said to be related by a **global phase**, and are indistinguishable by measurement:

$$\begin{aligned} |\langle e_i | \psi' \rangle|^2 &= |e^{i\Theta} \langle e_i | 4 \rangle|^2 \\ &= (e^{i\Theta} \langle e_i | 4 \rangle)(e^{-i\Theta} \langle e_i | 4 \rangle) \\ &= |\langle e_i | 4 \rangle|^2 \end{aligned}$$

On the other hand, a **relative phase**, p. 9.

$$|4\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle), \quad |4'\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$$

is distinguishable!

*relative phase  
on |1> state*

$$|\langle + | 4 \rangle|^2 = 1 \quad |\langle + | 4' \rangle|^2 = 0$$

$$|\langle - | 4 \rangle|^2 = 0 \quad |\langle - | 4' \rangle|^2 = 1$$

So given either  $|4\rangle$  or  $|4'\rangle$ , you can determine which you have by measuring in the  $\{|+\rangle, |-\rangle\}$  basis.

## (The Bloch sphere)

A final note about qubits is that we can visualize their states as lying on a **3-dimensional unit sphere** called the **Bloch sphere**.

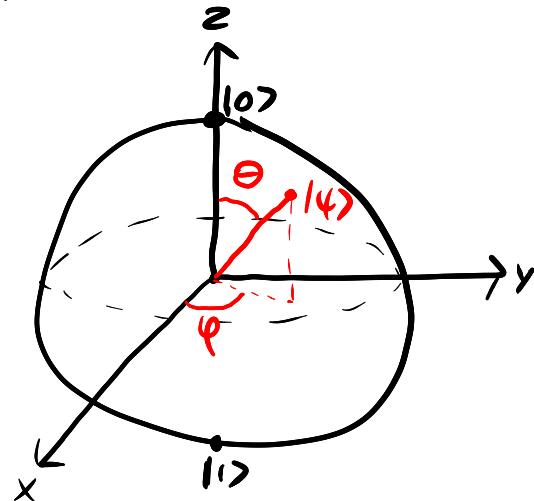
Since the state of a single qubit is

$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle, |\alpha|^2 + |\beta|^2 = 1$$

we may write it up to global phase as

$$|\psi\rangle = \cos(\theta/2)|0\rangle + e^{i\varphi}\sin(\theta/2)|1\rangle$$

These angles  $(\theta, \varphi)$  define a point on the unit 3-sphere like so



In this picture, we can view the **relative phase**  $\varphi$  as rotating our state around the  $z$ -axis. While the Bloch sphere has limited use in higher-dimensional or multi-qubit systems, it can be useful for understanding **single-qubit unitaries**, which as we will see later on correspond to **rotations** of the Bloch sphere.